

Forced hybrid-kinetic turbulence in 2D3V

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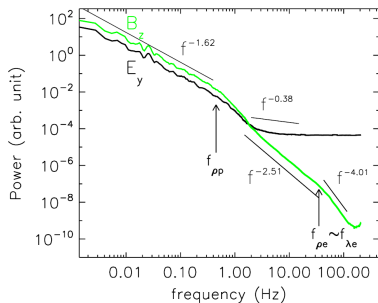
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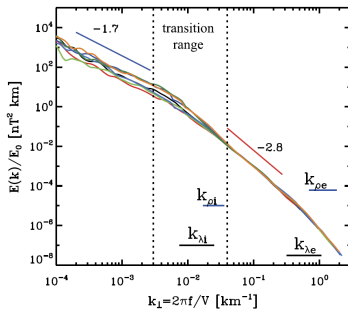
- 1 Motivation
- 2 The hybrid Vlasov-Maxwell (HVM) model
- 3 Results

Solar wind (SW) turbulence below the ion gyroradius

SW in-situ satellite measurements of turbulent energy spectra



[Sahraoui et al., PRL **102** (2009)]

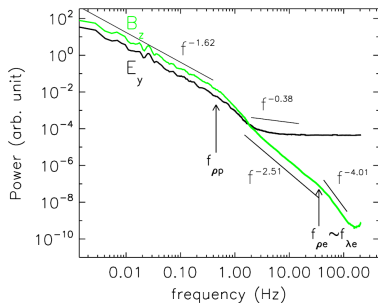


[Alexandrova et al., Space Sci Rev **178** (2013)]

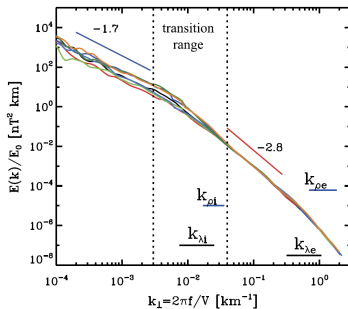
- **large scales:** magnetohydrodynamic (MHD) inertial range $\rightarrow \sim k_{\perp}^{-5/3}$ spectrum.
- **first spectral break** at ions' characteristic scales ($k_{\perp} \rho_i \sim 1$ and/or $k_{\perp} d_i \sim 1$).
- “dissipation/dispersion” range ($1 \lesssim k_{\perp} \rho_i \lesssim \rho_i/\rho_e$):
 - \rightarrow **B-field spectrum:** slope in the range $[-2.5, -3]$.
 - \rightarrow **E-field spectrum:** slope in the range $[-0.3, -1.3]$ (\rightarrow noise?).
 - \rightarrow energy in the E-field overcomes the magnetic counterpart.

Solar wind (SW) turbulence below the ion gyroradius

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[Alexandrova et al., Space Sci Rev **178** (2013)]

A long-lasting debate and open problem in SW turbulence research:

what is the **nature** of turbulent fluctuations below ion kinetic scales?

Solar wind (SW) turbulence below the ion gyroradius

Theoretical candidates:

kinetic Alfvén waves (KAWs)

[Schekochihin et al., ApJ Supp. Series **182**, 310 (2009)]

$$E_B(k_{\perp}) \propto k_{\perp}^{-7/3}$$

$$E_E(k_{\perp}) \propto k_{\perp}^{-1/3}$$

whistler waves

[Galtier & Bhattacharjee, PoP **10**, 3065 (2003)]

$$E_B(k_{\perp}) \propto k_{\perp}^{-7/3}$$

$$E_E(k_{\perp}) \propto k_{\perp}^{-1/3}$$

Same spectra, but different physics



auxiliary methods to distinguish between them

Possible sources of steepening:

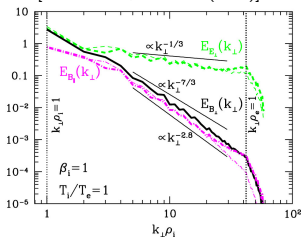
- Landau damping [Howes et al., JGR **113** (2008)]
- Compressibility effects: $E_B \propto k_{\perp}^{-7/3-2\xi}$ [Alexandrova et al., ApJ **674** (2008)]
- Intermittency corrections: $E_B \propto k_{\perp}^{-8/3}$ and $E_E \propto k_{\perp}^{-2/3}$ [Boldyrev & Perez, ApJL **758** (2012)]

Solar wind (SW) turbulence below the ion gyroradius

Numerical simulations: reproducing energy spectra

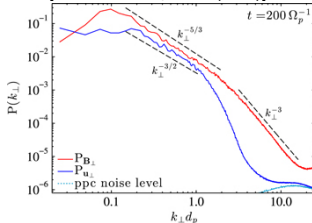
3D GK driven KAWs

[Howes et al., PRL 107 (2011)]



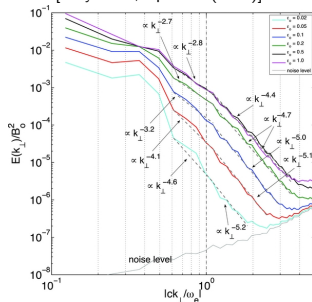
2D hybrid-PIC freely-decaying

[Franci et al., ApJ 812 (2015)]



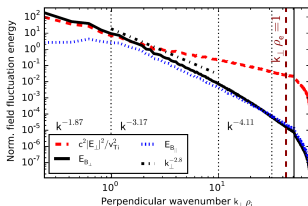
3D PIC freely-dec. whistlers

[Gary et al., ApJ 755 (2012)]



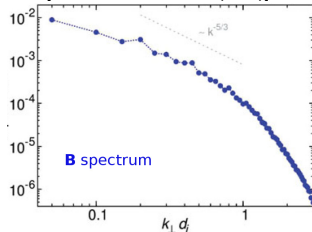
3D GK driven KAWs

[Told et al., PRL 115 (2015)]



3D HVM freely-decaying

[Servidio et al., JPP 81 (2015)]



So far, freely-decaying simulations and/or focus on one scenario at a time

Hybrid Vlasov-Maxwell (HVM) model

Fully kinetic ions & massless electron fluid:

[Valentini et al., JCP 225, 753 (2007)]

$$\frac{\partial f_i}{\partial t} + \mathbf{v} \cdot \frac{\partial f_i}{\partial \mathbf{x}} + (\mathbf{E} + \mathbf{v} \times \mathbf{B} + \mathbf{F}) \cdot \frac{\partial f_i}{\partial \mathbf{v}} = 0 \quad (\text{Vlasov equation})$$

$$\mathbf{E} = -\mathbf{u}_i \times \mathbf{B} + \frac{1}{n} (\mathbf{J} \times \mathbf{B}) - \frac{1}{n} \nabla P_e + \eta \mathbf{J} + \mathcal{O}\left(\frac{m_e}{m_i}\right) \quad (\text{gener. Ohm's law})$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}, \quad \nabla \times \mathbf{B} = \mathbf{J} + \frac{\partial \mathbf{E}}{\partial t} \quad (\text{Maxwell's equations})$$

$\mathbf{F} = \mathbf{F}(\mathbf{x}, t)$: random forcing, δ -correlated in time.

$$m_e = 0, \quad n_i = n_e = n, \quad \omega/k \ll c, \quad P_e = nT_{e0}$$

Simulations setup

- **2D-3V phase space:**

$$1024^2 \times 51^3 \text{ grid points } (\mathbf{k}_\perp d_i \in [0.1, 51.2])$$

- **initial condition:**

$$f_i(\mathbf{x}, \mathbf{v}, t = 0) = \text{isotropic Maxwellian}$$

$$\mathbf{B}(\mathbf{x}; t = 0) = B_0 \mathbf{e}_z + \delta \mathbf{B}(\mathbf{x}) \quad (|\delta \mathbf{B}| \ll B_0 \text{ and } 0.1 \leq (k_\perp d_i)_{\delta B} \leq 0.3)$$

- **F injection scale:**

$$0.1 \leq (k_\perp d_i)_F \leq 0.2 \text{ (continuously forced)}$$

→ forcing contributions: $\sim 50\%$ compressible, $\sim 50\%$ incompressible

- **beta regimes investigated:**

$$\beta = 0.2, 1 \text{ and } 5$$

The quest for a compromise: model & setup

Major “weak” points

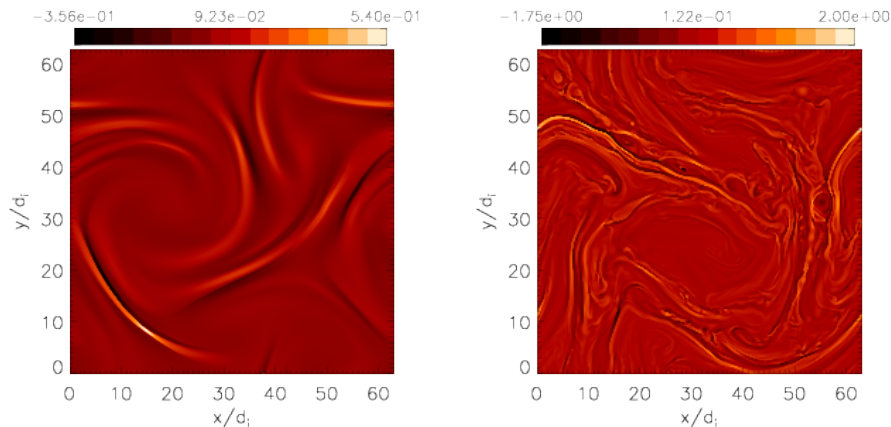
- reduced dimensionality (2D) of the simulations
- electron Landau damping (LD) is missing on all modes

Major “strong” points

- in 2D we can include three decades in the spectra
- fully kinetic ions (e.g., ion cyclotron resonances are included)
- we do not focus on a particular mode (both KAWs and whistler are allowed)
- **F** allows to reach a quasi-steady turbulent state
- the growth of in-plane magnetic fluctuations allows for $k_{\parallel} \neq 0$

we expect these “2.5D” simulations to retain some important dynamical features of the fully 3D case

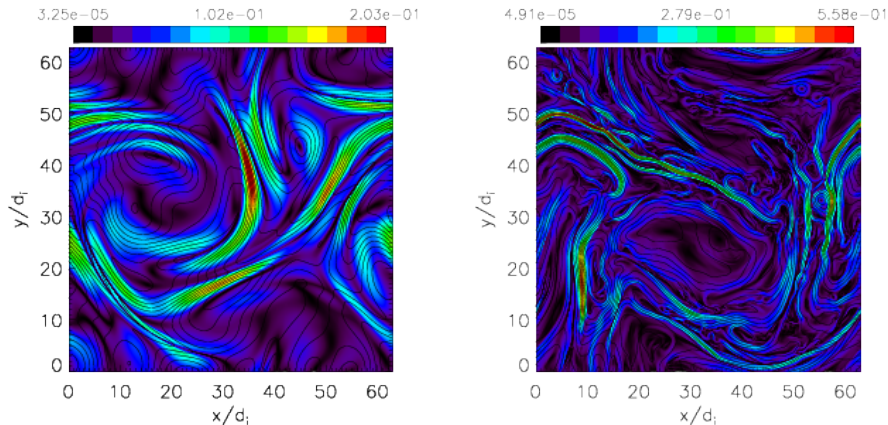
Developing plasma turbulence (J_z)



Example of J_z contours for $\beta_i = 1$, at $\Omega_{ci} t = 120$ (left) and $\Omega_{ci} t = 225$ (right).

- formation of **small-scale structures** \rightarrow **kinetic regime**
- current sheets \rightarrow magnetic reconnection \rightarrow fully developed turbulence

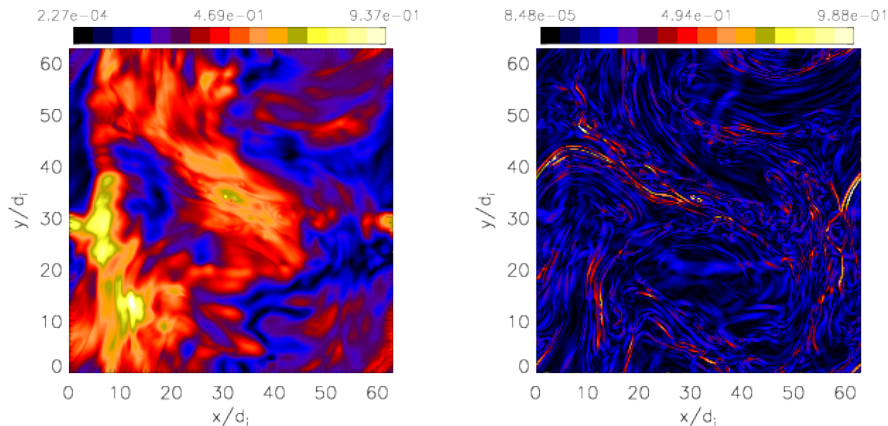
Developing plasma turbulence (B_{\perp})



Example of B_{\perp} contours and A_z lines for $\beta_i = 1$, at $\Omega_{ci}t = 120$ (left) and $\Omega_{ci}t = 225$ (right).

- in-plane magnetic fluctuations: randomly oriented, $\langle B_{\perp} \rangle < 0.1$
- local high- B_{\perp} spots: **current sheets**, **coherent structures**

Developed plasma turbulence (E_{\perp})

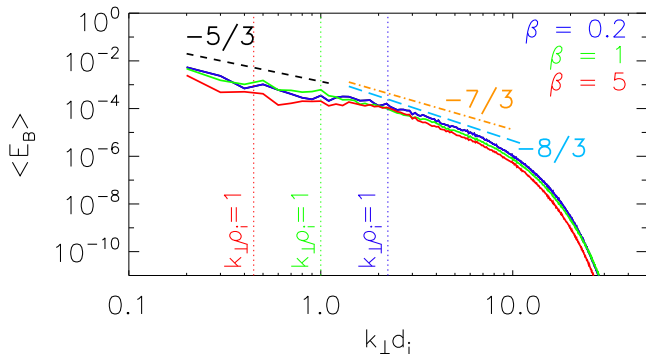


Contours of E_{MHD} (left) and of E_{Hall} (right) for $\beta_i = 1$ at $\Omega_{ci} t = 225$.

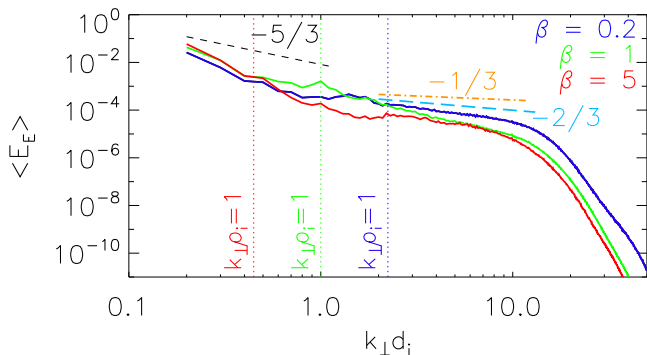
- $\mathbf{E}_{\text{MHD}} = \mathbf{u}_i \times \mathbf{B}$ dominates at **large-scales** (left)
- $\mathbf{E}_{\text{Hall}} = (\mathbf{J} \times \mathbf{B})/n$ dominates at **small-scales**, inside current sheets (right)

Magnetic energy spectrum

Cerri et al., ApJL 822, L12 (2016)



- $k_{\perp} d_i < 1$: Kolmogorov-type $k_{\perp}^{-5/3}$ spectrum
- spectral break at $1 \lesssim k_{\perp} d_i \lesssim 2$
- $k_{\perp} d_i > 1$: consistent with $k_{\perp}^{-8/3}$ at $\beta = 0.2, 1$ (k_{\perp}^{-3} at $\beta_i = 5$)



- electric energy overcomes magnetic counterpart at $k_{\perp} d_i \sim 2$
- spectral slopes generally steeper than theory predictions (observed in other simulations and some SW measurements \rightarrow feedbacks?)

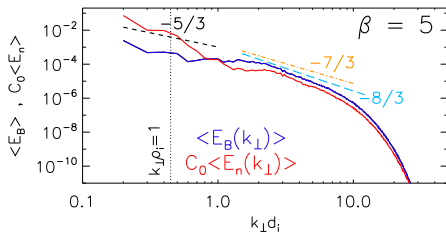
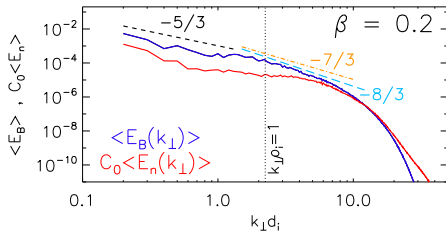
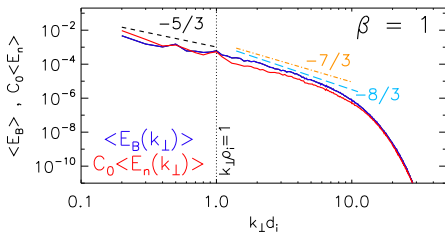
KAWs or whistlers? (Auxiliary method I)

Auxiliary method I:

[Chen et al., PRL **110**, 225002 (2013)]

comparing the level of E_B and $C_0 E_n$
(with $C_0 = [\beta_i(1 + \tau)/2][1 + \beta_i(1 + \tau)/2]$)

- **KAWs** $\rightarrow C_0 E_n \simeq E_B$.
- **whistlers** $\rightarrow C_0 E_n \ll E_B$.



Cerri et al., ApJL **822**, L12 (2016)

KAWs or whistlers? (Auxiliary method II)

Auxiliary method II:

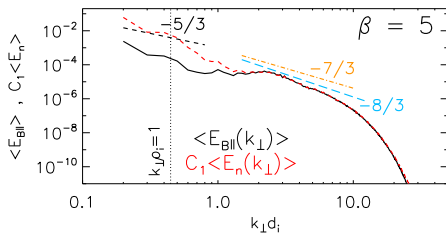
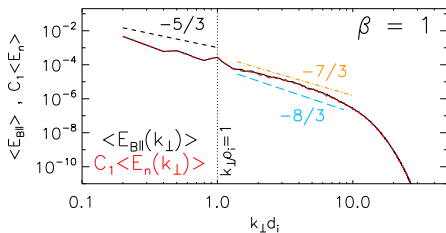
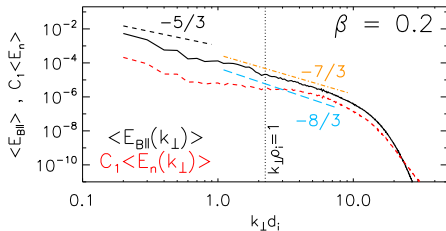
[Boldyrev et al., ApJ 777, 41 (2013)]

KAWs fluctuations would obey the following relation:

$$C_1 E_n = E_{B\parallel}$$

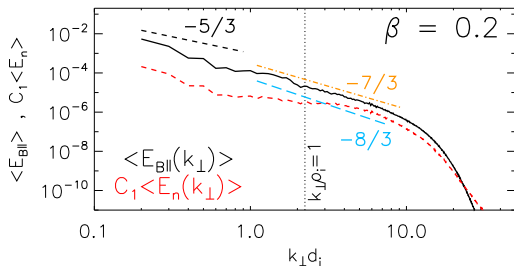
(with $C_1 = [\beta_i(1 + \tau)/2]^2$)

Cerri et al., ApJL 822, L12 (2016)



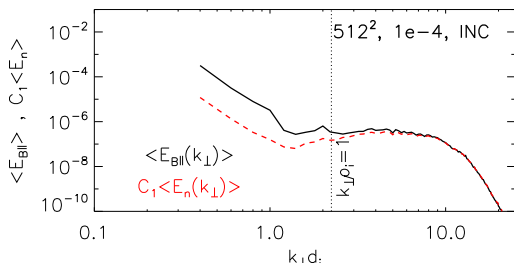
Partially compressible vs incompressible injection ($\beta = 0.2$)

Partially compressible forcing ($\nabla \cdot \mathbf{F} \neq 0$):



→ well separated even
at $k_{\perp} \rho_i > 1$

Purely incompressible forcing ($\nabla \cdot \mathbf{F} = 0$):



→ transition to KAWs
at $k_{\perp} \rho_i \sim 1$

- **general agreement of spectral properties** of the turbulence (e.g., power laws and spectral breaks) with observations/theory.
- in this setup turbulence mainly involves **whistler fluctuations at low β** , and **KAWs at somewhat higher β** .
- KAW \leftrightarrow whistler **turbulence transition**: possible correlation with **resonant/non-resonant damping** of the modes.
(not straightforward: linear damping and/vs non-linear effects)
- **compressibility level** of injected fluctuations matters \rightarrow non-universality and possible implications on time and space variability of SW.

 \rightarrow **call for further investigations on these topics...**

Thanks for your attention!